

# LONG TIME EXISTENCE OF HYPERBOLIC RICCI-BOURGUIGNON FLOW ON RIEMANNIAN SURFACES

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**Abstract.** We consider the *hyperbolic Ricci-Bourguignon flow (HRBF)* equation on Riemannian surfaces and we find a sufficient and necessary condition to this flow has global classical solution. Also, we show that the scalar curvature of the solution metric  $g_{ij}$  convergence to the flat curvature.

*Key words and Phrases:* Hyperbolic geometric flow; quasilinear hyperbolic equation; strict hyperbolicity

## 1. INTRODUCTION

Let  $(M, g)$  be an  $n$ -dimensional complete Riemannian manifold with Riemannian metric  $g_{ij}$ . The general variation equation

$$\frac{\partial^2 g_{ij}}{\partial t^2} + 2R_{ij} + \mathcal{F}(g, \frac{\partial g}{\partial t}) = 0, \quad (1)$$

was introduced by Kong and Liu ([4]) and called the generalized *hyperbolic geometric flow* (denoted by HGF). Here  $\mathcal{F}$  are some smooth functions of the Riemannian metric and its first derivative with respect to  $t$ , and we consider  $R_{ij}$  as the components of Ricci curvature tensor. Liu and Zhang in ([8]) have shown that the *hyperbolic geometric flow* (HGF) has global classical solution on Riemannian surfaces. In this paper, we would like to prove that the global solution of hyperbolic Ricci-Bourguignon flow (HRBF) exists on Riemannian surfaces.

The present work investigates the variation of a Riemannian metric  $g_{ij}$  on a Riemannian surface  $M$  by its Ricci curvature tensor  $R_{ij}$  and scalar curvature  $R$  under

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the following equation

$$\frac{\partial^2 g_{ij}}{\partial t^2} = -2R_{ij} + 2\rho Rg_{ij} \tag{2}$$

where  $\rho$  is a real constant. When  $\rho = 0$ , this equation is hyperbolic geometric flow and the global existence and blowup phenomenon of smooth solutions to this flow on Riemannian surface have been investigated in [8]. The Ricci-Bourguignon flow is  $\frac{\partial g_{ij}}{\partial t} = -2R_{ij} + 2\rho Rg_{ij}$  and the short time existence and uniqueness for solution to the Ricci-Bourguignon flow on  $[0, T)$  were showed by Catino et al ([1]) for  $\rho < \frac{1}{2(n-1)}$ .

This study regards the initial metric as follows

$$ds^2 = u_0(x)(dx^2 + dy^2) \quad \text{at } t = 0 \tag{3}$$

on a surface of topological type  $\mathbb{R}^2$ , where  $u_0(x)$  is a function from  $C^2$  class with bounded  $C^2$  norm and the following inequality is hold

$$0 < k \leq u_0(x) \leq m < \infty \tag{4}$$

where  $k$  and  $m$  are positive constants.

Since all the information about curvature is contained in the scalar curvature function  $R$ , we can simplify the HRBF equation on this surface. According to our notation,  $R = 2K$ , where  $K$  denotes Gauss curvature and also the Ricci curvature is given by  $R_{ij} = \frac{1}{2}Rg_{ij}$ , so the (HRBF) equation simplifies to

$$\frac{\partial^2 g_{ij}}{\partial t^2} = -R(1 - 2\rho)g_{ij}. \tag{5}$$

At least locally the metric for a surface can be written as  $g_{ij} = u(t, x, y)\delta_{ij}$ , where  $u(t, x, y) > 0$ , and  $\delta_{ij}$  is Kronecker's symbol. Hence, we have

$$R = -\frac{\Delta \ln u}{u} \tag{6}$$

as a result, the aforementioned equation (5) reduces to  $u_{tt} - (1 - 2\rho)\Delta \ln u = 0$ . The initial data  $u_0(x)$  depends only on  $x$  and not  $y$ ; thus, we can consider the Cauchy problem as below

$$\begin{cases} u_{tt} - (1 - 2\rho)(\ln u)_{xx} = 0, \\ u = u_0(x), & \text{at } t = 0 \\ u_t = u_1(x), & \text{at } t = 0 \end{cases} \tag{7}$$

where  $u_1(x) \in C^1$  with bounded  $C^1$  norm. By using the transformation

$$\phi = \ln u, \tag{8}$$

Kong and Liu in ([5]) proved a theorem as follows

**Theorem 1.1.** *Suppose that  $u_1(x) \geq |u'_0(x)|/\sqrt{u_0(x)}$  for all  $x \in \mathbb{R}$ . Then, the Cauchy problem (7) admits a unique global solution for all  $t \in \mathbb{R}$ .*

*Moreover, if  $u_1(x) \equiv u'_0(x)/\sqrt{u_0(x)}$ , and there exists a point  $x_0 \in \mathbb{R}$  such*

that  $u'_0(x_0) < 0$ , then the Cauchy problem (7) admits a unique classical solution only in  $[0, T) \times \mathbb{R}$ , where

$$T = -\frac{2}{\inf_x \{u'_0(x)u_0^{-3/2}(x)\}}. \quad (9)$$

The following theorem will be proven without using (6) in our investigation.

**Theorem 1.2.** *Let*

$$u_1(x) + \frac{u'_0(x)}{\sqrt{u_0(x)}} \geq 0 \quad \text{at all } x \in \mathbb{R}, \quad (10)$$

and

$$u_1(x) - \frac{u'_0(x)}{\sqrt{u_0(x)}} \geq 0 \quad \text{at all } x \in \mathbb{R}. \quad (11)$$

Hence, the Cauchy problem (7) has a unique global solution for all  $t \in \mathbb{R}$ .

**Theorem 1.3.** *If a point  $x_0 \in \mathbb{R}$  exists, which satisfy*

$$u_1(x_0) + \frac{u'_0(x_0)}{\sqrt{u_0(x_0)}} < 0 \quad (12)$$

or there exists a point  $x_0 \in \mathbb{R}$ , such that

$$u_1(x_0) - \frac{u'_0(x_0)}{\sqrt{u_0(x_0)}} < 0 \quad (13)$$

thus, the Cauchy problem (7) has a unique classical solution only in  $[0, T) \times \mathbb{R}$ .

**Note.** Based on Theorem 1.2, we can conclude the Cauchy problem

$$\begin{cases} \frac{\partial^2 g_{ij}}{\partial t^2} = -2R_{ij} + 2\rho R g_{ij}, & \text{for } i, j = 1, 2 \\ g_{ij} = u_0(x)\delta_{ij}, & \text{for } t = 0 \\ \frac{\partial g_{ij}}{\partial t} = u_1(x)\delta_{ij}, & \text{for } t = 0 \text{ and } i, j = 1, 2. \end{cases}$$

has a unique smooth solution for all  $t \in \mathbb{R}$ . Besides we can consider the solution metric  $g_{ij}$  as below

$$g_{ij} = u(x, t)\delta_{ij} \quad \text{for } i, j = 1, 2. \quad (14)$$

We will prove the above mentioned Theorems 1.2 and 1.3, in the subsequent sections (3 and 4, respectively). Moreover, using Theorem 1.2, the following theorem will be proven in Section 5.

**Theorem 1.4.** *Let be*

$$\inf_x \left\{ u_1(x) + \frac{u'_0(x)}{\sqrt{u_0(x)}} \right\} > 0 \quad \text{and} \quad \inf_x \left\{ u_1(x) - \frac{u'_0(x)}{\sqrt{u_0(x)}} \right\} > 0. \quad (15)$$

Hence, a unique classical solution of (1) exists as the form (14) for all time. Furthermore, the scalar curvature  $R(x, t)$  relates to the solution metric  $g_{ij}$  admits

$$R(x, t) \rightarrow 0 \quad \text{as } t \rightarrow +\infty,$$

and  $R(x, t) \leq k_1$  for all  $(t, x) \in \mathbb{R}^+ \times \mathbb{R}$ , where  $k_1$  is a positive constant and independent of  $t$  and  $x$ .

2. PRELIMINARIES

In this section we require only to discuss the classical solution on  $t \geq 0$ . The result for  $t \leq 0$  can be easily obtained.

Suppose that

$$u_t = v \quad \text{and} \quad u_x = w. \tag{16}$$

Thus, from the above equations and Cauchy problem (7) we have

$$u_t = v, \quad w_t - v_x = 0, \quad \text{and} \quad v_t - \left(\frac{1-2\rho}{u}\right)w_x = (2\rho-1)\frac{w^2}{u^2}. \tag{17}$$

Eigenvalues of equations (17) can be easily calculated as follows

$$\lambda_1 = -\lambda, \quad \lambda_2 = 0, \quad \lambda_3 = \lambda, \quad \lambda = \sqrt{\frac{1-2\rho}{u}} \tag{18}$$

and we have the matrices  $L(U)$  and  $R(U)$ (where  $U = (u, w, v)$ ) of left and right eigenvectors, respectively as below

$$L(U) = \begin{pmatrix} l_1(U) \\ l_2(U) \\ l_3(U) \end{pmatrix} = \begin{pmatrix} 0 & \lambda & 1 \\ 1 & 0 & 0 \\ 0 & -\lambda & 1 \end{pmatrix},$$

$$R(U) = (r_1(u), r_2(u), r_3(u)) = \begin{pmatrix} 0 & 1 & 0 \\ \lambda & 0 & -\lambda \\ 1 & 0 & 1 \end{pmatrix}.$$

Equation system (17) is a linear degenerate strict hyperbolic system because of  $\nabla \lambda_i(U)r_i(U) \equiv 0$  for  $i = 1, 2, 3$ .

Define  $p$  and  $q$  as follows

$$p = v + \lambda w \quad \text{and} \quad q = v - \lambda w. \tag{19}$$

**Lemma 2.1.**  $p$  and  $q$  satisfy the following equations:

$$p_t - \lambda p_x = \frac{1}{4(1-2\rho)}\lambda^2(q-p)p, \tag{20}$$

$$u_t = \frac{1}{2}(p+q), \tag{21}$$

$$q_t + \lambda q_x = \frac{1}{4(1-2\rho)}\lambda^2(p-q)q. \tag{22}$$

*Proof.* By differentiating of the function  $\lambda$  with respect to  $t$  and  $x$ ,  $\lambda_t$  and  $\lambda_x$  can easily be obtained as,

$$\lambda_t = -\frac{1}{2(1-\rho)}\lambda^3 v,$$

and

$$\lambda_x = -\frac{1}{2(1-2\rho)}\lambda^3 w.$$

Therefore,

$$\begin{aligned} p_t - \lambda p_x &= (v + \lambda w)_t - \lambda(v + \lambda w)_x \\ &= v_t - \lambda v_x + \lambda(w_t - \lambda w_x) + w(\lambda_t - \lambda \lambda_x) \\ &= v_t - \lambda^2 w_x - \lambda(v_x - w_t) + w(\lambda_t - \lambda \lambda_x) \\ &= -\frac{1}{2(1-2\rho)} \lambda^3 w(v + \lambda w) = \frac{1}{4(1-2\rho)} \lambda^2 (q-p)p. \end{aligned}$$

We can prove (22) in the same way as above, and it is obvious that (21) is hold.  $\square$

For the next lemma, consider

$$r = p_x + \frac{1}{8(1-2\rho)} \lambda p q \quad \text{and} \quad s = q_x - \frac{1}{8(1-2\rho)} \lambda p q.$$

**Lemma 2.2.** *r and s satisfy*

$$r_t - \lambda r_x = \frac{\lambda^2}{4(1-2\rho)} (2q-3p)r + \frac{\lambda^3}{32(1-2\rho)} (2p-3q)pq + \frac{\lambda^3 p(q-p)}{32(1-2\rho)^2} (p+5q), \quad (23)$$

$$s_t + \lambda s_x = \frac{\lambda^2}{4(1-2\rho)} (2q-3p)r + \frac{\lambda^3}{32(1-2\rho)} (2p-3q)pq + \frac{\lambda^3 p(q-p)}{32(1-2\rho)^2} (p+5q). \quad (24)$$

*Proof.* Suppose

$$L_1 = \frac{\partial}{\partial t} - \lambda \frac{\partial}{\partial x} \quad \text{and} \quad L_2 = \frac{\partial}{\partial t} + \lambda \frac{\partial}{\partial x}.$$

Hence, by a direct computation we can get

$$\begin{aligned} L_1 p_x &= \frac{1}{4(1-2\rho)} \lambda^2 ((2q-3p)p_x + pq_x) + \frac{1}{8(1-2\rho)^2} \lambda^3 p(q-p)^2, \\ L_2 q_x &= \frac{1}{4(1-2\rho)} \lambda^2 ((2p-3q)q_x + qp_x) - \frac{1}{8(1-2\rho)^2} \lambda^3 q(p-q)^2. \end{aligned}$$

Now we can easily prove (23) and (24).  $\square$

Notice from Cauchy problem (7), (16) and (19), we can write following equations at  $t = 0$ .

$$p = p_0(x) \equiv u_1(x) + \lambda_0(x)u'_0(x), \quad u = u_0(x) \quad (25)$$

$$q = q_0(x) \equiv u_1(x) - \lambda_0(x)u'_0(x), \quad (26)$$

where  $\lambda_0(x) = \sqrt{\frac{1-2\rho}{u_0(x)}}$ . Now in following theorem we show that the Cauchy problem (7) has a unique global solution under some conditions.

**Theorem 2.3.** *Let  $M_1$  be a positive constant satisfying*

$$0 \leq p(x, t) \leq M_1 \quad \text{and} \quad 0 \leq q(x, t) \leq M_1, \quad (27)$$

*then, on  $D(T)$ ,*

$$|u(x, t)| \leq M(T), \quad |u_x(x, t)| \leq M(T), \quad |u_t(x, t)| \leq M(T),$$

$$|r(x, t)| \leq M(T), \quad |s(x, t)| \leq M(T),$$

where  $M(T)$  is a positive constant, and

$$D(T) = \{(x, t) | x \in \mathbb{R}, 0 \leq t \leq T, T > 0\}.$$

Hence, the Cauchy problem (7) has a unique global classical solution on  $t \geq 0$ , by the local existence theorem of the classical solution to quasilinear hyperbolic systems.

*Proof.* Consider any point  $(t, x)$ , and let

$$x = x_1(t, \beta_1), \quad x = x_2(t, \beta_2), \quad x = x_3(t, \beta_3)$$

be the  $\lambda_1, \lambda_2$  and  $\lambda_3$  characteristics, respectively, that satisfy

$$(x_1)_t = \lambda_1 = -\lambda, \quad (x_2)_t = \lambda_2 = 0, \quad (x_3)_t = \lambda_3 = \lambda$$

$$x_1(0, \beta_1) = \beta_1, \quad x_2(0, \beta_2) = \beta_2, \quad x_3(0, \beta_3) = \beta_3.$$

By integrating (21) along the  $\lambda_2$  characteristics we can obtain

$$u(x, t) = u_0(\beta_2) + \frac{1}{2} \int_0^t (p + q)(x_2(\tau, \beta_2), \tau) d\tau. \tag{28}$$

Thus, as a result of (21), (27) and (28) we have

$$|u_t| \leq M_1(T) \quad \text{and} \quad 0 < \inf_x u_0(x) \leq u(x, t) \leq M_2(T).$$

Using a same method, by integrating (23) along the  $\lambda_1$  characteristics  $x = x_1(t, \beta_1)$

$$|r(x, t)| \leq M_2(T) + M_3(T) \int_0^t R(\tau) d\tau, \quad \text{where } R(t) = \sup_x |r(x, t)|.$$

Therefore, we have  $|r(x, t)| \leq M_4(T)$  by the Bellman lemma. As a similar way, holds  $|s(x, t)| \leq M_5(T)$ . Since  $(u_x)_t = \frac{1}{2}(r + s)$ , it is obvious that

$$|u_x(x, t)| \leq M_6(T),$$

which  $M_i(t)$  for  $i = 1, 2, 3, \dots$  denote positive constants. □

### 3. PROOF OF THEOREM 1.2

On the basis of the local existence and uniqueness theorems of the classical solutions to the quasilinear hyperbolic systems ([7]), to prove Theorem (1.2) it suffices to establish uniform *a priori* estimates of the  $C^1$  norms of  $p, q$  and  $u$ . We have following lemma from [2, 3].

**Lemma 3.1.** *Suppose*

$$u_t + \lambda_1(x, t)u_x = A(x, t)(u - v),$$

$$v_t + \lambda_2(x, t)v_x = B(x, t)(v - u),$$

where  $\lambda_1, \lambda_2, A$  and  $B$  are continuous functions, and  $\lambda_1 \leq \lambda_2$ . If  $A$  and  $B$  are both non positive, then

$$\min(u_0(x), v_0(x)) \leq u(x, t), v(x, t) \leq \max(u_0(x), v_0(x)).$$

For prove Theorem 1.2 we need the following lemma.

**Lemma 3.2.** *On the existence domain of the classical solution to the Cauchy problem (7) and (25), if (10) and (11) hold, then*

$$0 \leq p(x, t) \leq \sup_{x \in \mathbb{R}} p_0(x), \quad (29)$$

$$0 \leq q(x, t) \leq \sup_{x \in \mathbb{R}} q_0(x), \quad (30)$$

$$0 < \inf_{x \in \mathbb{R}} u_0(x) \leq u(x, t) \leq \sup_x u_0(x) + Ct, \quad (31)$$

where  $C > 0$  is a constant.

*Proof.* Along  $\lambda_1$  characteristics, we can obtain

$$p(x, t) = p_0(\beta_1) \exp\left(\int_0^t \frac{1}{4(1-2\rho)} \lambda^2 (q-p)(x_1(\tau, \beta_1), \tau) d\tau\right).$$

By (25) and (10),  $p_0(x) \geq 0$  for all  $x \in \mathbb{R}$ . Therefore, we have  $p(x, t) \geq 0$ . In a similar way we can prove  $q(x, t) \geq 0$ . As a result of these inequalities, we have

$$\frac{1}{4(1-2\rho)} \lambda^2 p \geq 0 \quad \text{and} \quad \frac{1}{4(1-2\rho)} \lambda^2 q \geq 0.$$

Hence, by Lemma 3.1 we can easily see that

$$p(x, t) \leq \sup_x p_0(x) \quad \text{and} \quad q(x, t) \leq \sup_x q_0(x).$$

Also, we can get following equality by integrating (21)

$$u(x, t) = u_0(\beta_2) + \frac{1}{2} \int_0^t (p+q)(x_2(\tau, \beta_2), \tau) d\tau.$$

Thus, we can get to result since  $p(x, t) \geq 0$  and  $q(x, t) \geq 0$ .  $\square$

*Proof of Theorem 1.2.* Now from aforementioned Lemma 3.2 and Theorem 2.3, Theorem 1.2 is obvious.  $\square$

**Note.** By (6) and (29), and based on the hypotheses of Theorem 1.2, we have

$$|R(x, t)| \leq M_7(T).$$

#### 4. PROOF OF THEOREM 1.3

The blow-up phenomena of the *hyperbolic geometric flow* will be discussed in this section.

Suppose

$$m = \sqrt{\lambda} p \quad \text{and} \quad n = \sqrt{\lambda} q. \quad (32)$$

We have

$$\frac{1}{4} \lambda^2 q = \frac{1-2\rho}{4} ((\ln u)_t - \lambda (\ln u)_x) \quad \text{and} \quad \frac{1}{4} \lambda^2 p = \frac{1-2\rho}{4} ((\ln u)_t + \lambda (\ln u)_x).$$

By the use of (16), (19) and Lemma 2.1 the following lemma can be proven

**Lemma 4.1.** *m and n satisfy*

$$m_t - \lambda m_x = -\frac{1}{4(1-2\rho)}\lambda^{3/2}m^2, \tag{33}$$

$$n_t - \lambda n_x = -\frac{1}{4(1-2\rho)}\lambda^{3/2}n^2. \tag{34}$$

Observe that at  $t = 0$ , set

$$m = m_0(x) = \sqrt[4]{\frac{1-2\rho}{u_0(x)}}\left(u_1(x) + \frac{u'_0(x)}{\sqrt{u_0(x)}}\right), \tag{35}$$

$$n = n_0(x) = \sqrt[4]{\frac{1-2\rho}{u_0(x)}}\left(u_1(x) - \frac{u'_0(x)}{\sqrt{u_0(x)}}\right). \tag{36}$$

*Proof of Theorem 1.3.* Without loss of generality, we assume that (12) holds; in the same way we can proceed if (13) holds.

As a result of (33) and (34) we have  $m_t - \lambda m_x \leq 0$  and  $n_t - \lambda n_x \leq 0$ . Thus, we can easily see that

$$m(x, t) + n(x, t) \leq M_0 \quad \text{and} \quad M_0 \equiv \sup m_0(x) + \sup n_0(x). \tag{37}$$

Notice that  $u_0(x) \geq k > 0$ , and also from (12) and (35) we have  $m_0(x) < 0$ . Next, the get following equation is obtained from (33) by integrating along  $\lambda_1$  characteristics. That is,

$$m(x_0, t) = m_0(x_0)/F(t, x_0), \tag{38}$$

where

$$F(t, x_0) = 1 + m_0(x_0)/4(1-2\rho) \int_0^t \lambda^{3/2}(x_1(x_0, \tau), \tau) d\tau \text{ and } \lambda^{3/2} = \left(\frac{u}{1-2\rho}\right)^{-3/4}. \tag{39}$$

By (21) and (32), it is easy to see that  $\left(\left(\frac{u}{1-2\rho}\right)^{-3/4}\right)_t = \frac{3}{8(1-2\rho)}(m+n)$ . Hence, we have

$$\left(\frac{u}{1-2\rho}\right)^{3/4}(x, t) = \left(\frac{u_0}{1-2\rho}\right)^{3/4}(x_0) + \frac{3}{8(1-2\rho)} \int_0^t (m+n)(x_2(x_0, \tau), \tau) d\tau. \tag{40}$$

By (4), (37) and (40), we get

$$u^{3/4}(x, 0) \geq k^{3/4} \quad \text{and} \quad u^{3/4}(x, t) \geq M^{3/4} + \frac{3}{8(1-2\rho)}M_0t. \tag{41}$$

We consider three cases.

*Case(i).* If  $M_0 < 0$ , then there exists  $\tau_0 = 8(1-2\rho)M^{3/4}/(3(-M_0)) > 0$ , such that

$$u(x, t) \leq 0 \quad \text{and} \quad t \geq \tau_0.$$

This imply the system in (7) is meaningless for  $t \geq \tau_0$ , that is, it admits a unique local classical solution.



*Case(ii).* If  $M_0 = 0$ , then, by (39) and (41), following inequality can be easily obtained

$$F(x_0, t) \leq 1 + \frac{1-2\rho}{4} m_0(x_0) M^{-3/4} t.$$

When  $F(x_0, 0) = 1 > 0$  and  $m_0(x_0) < 0$ , we can find  $t_0 = 4M^{3/4}/(1-2\rho)(-m_0(x_0)) > 0$ , such that

$$F(x_0, t) \rightarrow 0^+ \quad \text{as } t \rightarrow t_0^-.$$
 (42)

Thus, the finite time  $T = T(x_0) > 0$  exists such that

$$F(x_0, t) \rightarrow -\infty \quad \text{as } t \rightarrow T^-.$$
 (43)

*Case(iii).* If  $M_0 > 0$ , then, by (39) and (41) we can get

$$F(x_0, t) \leq 1 + \frac{2m_0(x_0)}{3M_0} \ln\left(1 + \frac{3M_0}{8(1-2\rho)M^{3/4}} t\right).$$

Therefore, since,  $F(x_0, 0) = 1 > 0$  and  $m_0(x_0) < 0$ , there exists  $t_* > 0$  such that  $F(x_0, t) \rightarrow 0^+$  as  $t \rightarrow t_*^-$ , and then (43) follows.  $\square$

## 5. PROOF OF THEOREM 1.4

In this section, we will study the asymptotic behaviour of the scalar curvature  $R(x, t)$ .

*Proof of Theorem 1.4.* We assume that (15) holds. Using Lemmas 3.1 and 3.2, we have

$$K_1 \leq p(x, t) \leq K_2 \quad \text{and} \quad K_1 \leq q(x, t) \leq K_2,$$

where here after  $C_i$  for  $i = 1, 2, \dots$  denote positive constants independent of  $t$  and  $x$ . Therefore, we can obtain the following inequality as a result of (4) and (28)

$$K_3(1+t) \leq u(x, t) \leq K_4(1+t),$$
 (44)

and then

$$\frac{K_3^{3/4}(1+t)^{3/4}}{1-2\rho} \leq \frac{u^{3/4}(x, t)}{1-2\rho} \leq \frac{K_4^{3/4}(1+t)^{3/4}}{1-2\rho} \leq \frac{K_4^{3/4}(1+t)}{1-2\rho}.$$

It follows from(40) and (44) that

$$K_5(1 + \ln(1+t)) \leq F(x, t) \leq K_6(1 + (1+t)^{1/4}).$$

Thus, by (15) and (38), we get

$$0 \leq \frac{K_7}{1 + (1+t)^{1/4}} \leq m(x, t) \leq \frac{K_8}{1 + \ln(1+t)}.$$
 (45)

Similarly, we can obtain

$$0 \leq \frac{K_7}{1 + (1+t)^{1/4}} \leq m(x, t) \leq \frac{K_8}{1 + \ln(1+t)}.$$
 (46)

Hence,  $m(x, t) \rightarrow 0$  and  $n(x, t) \rightarrow 0$  as  $t \rightarrow +\infty$ . Noting (44), (45) and (46), we get

$$p = \frac{1}{\sqrt{\lambda}} m = \frac{u}{1-2\rho} m, \quad q = \frac{1}{\sqrt{\lambda}} n = \frac{u}{1-2\rho} n, \quad u_x = \frac{p-q}{2\lambda} = \frac{1}{2} u^{3/4} (m-n),$$

and

$$-\frac{K_8}{1 + \ln(1 + t)} \leq m(x, t) - n(x, t) \leq \frac{K_8}{1 + \ln(1 + t)}. \tag{47}$$

Then, we can easily obtain

$$|u_x| \leq K_9 \frac{(1 + t)^{3/4}}{1 + \ln(1 + t)}. \tag{48}$$

Next an easy derivation gives

$$u_{xx} = \frac{p_x - q_x}{2\lambda} + \frac{1}{2}\lambda^2 u_x^2 = \frac{1}{2}u^{1/2}(p_x - q_x) + \frac{1}{2u}u_x^2. \tag{49}$$

Let  $\bar{p} = p u$  and  $\bar{q} = q u$ . Thus, based on [5] we have

$$\begin{aligned} L_1 \bar{p}_x &= -A_1 \bar{p}_x - B_1 \bar{q}_x, \\ L_2 \bar{q}_x &= -A_2 \bar{p}_x - B_2 \bar{q}_x, \end{aligned}$$

where

$$A_1 = \frac{1}{4}(2\bar{q} + 3\bar{q}), \quad B_1 = \frac{3}{4}\bar{q}, \quad A_2 = \frac{3}{4}\bar{q}, \quad B_2 = \frac{1}{4}(2\bar{p} + 3\bar{q}).$$

Therefore, by [5] holds

$$|\bar{p}_x(x, t)|, |\bar{q}_x(x, t)| \leq K_{10}. \tag{50}$$

Noting that  $p_x - q_x = u_x(p - q)/u + u(\bar{p}_x - \bar{q}_x)$ , and from (6) and (49), we have

$$R = \frac{1}{u^3}(u_x^2 - uu_{xx}) = \frac{u_x^2}{2u^3} - \frac{p_x - q_x}{2u^{3/2}} = \frac{u_x^2}{2u^3} - \frac{\bar{p}_x - \bar{q}_x}{2u^{1/2}} - \frac{u_x(m - n)}{2u^3}. \tag{51}$$

Thus, from (41), (47), (48), (50) and (51) we conclude that

$$|R(x, t)| \leq \frac{K_{11}}{(1 + \ln(1 + t))^2(1 + t)^{3/2}} + \frac{K_{12}}{(1 + t)^{1/2}} + \frac{K_{13}(1 + t)^{3/4}}{(1 + t)^3(1 + \ln(1 + t))^2}.$$

Hence,  $R(x, t) \rightarrow 0$  as  $t \rightarrow +\infty$ . □

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### REFERENCES

- [1] Catino, G., Cremaschi, Djadli, L.Z., Mantegazza C., and Mazzieri, L., "The Ricci-Bourguignon flow", *Pacific Journal of Mathematics*, **287**(2)(2017), 337-370.
- [2] Hong, J.X., "The global smooth solutions of Cauchy problems for hyperbolic equation of Monge-Ampère type", *Nonlinear Anal.* **24**(1995),1649-1663.
- [3] Kong, D.X., "Maximum principle in nonlinear hyperbolic systems and its applications", *Nonlinear Anal.* **32**(7)(1998), 871-880.
- [4] Kong D.X. and Liu, K., "Wave character of metrics and hyperbolic geometric flow", *J. Math. Phys.* **48**(10)(2007),103508.
- [5] Kong, D.X., Liu K., and Xu, D.L., "The hyperbolic geometric flow on Reimann surfaces", *Comm. Partial Differential Equations*, **34**(4-6)(2009), 553-580.
- [6] Li T.T. and Liu, F.G., "Singularity caused by eigenvalues for quasilinear hyperbolic systems", *Comm. Partial Differential Equations*, **28**(3-4)(2003),477-503.

- [7] Li T.T. and Yu, W.C., *Boundary value problems for quasilinear hyperbolic systems*, Duke University Math. Series 5, Duke University Press, Durham, NC, 1985.
- [8] Liu F. and Zhang, Y., "Global classical solutions to hyperbolic geometric flow on Riemann surfaces", *Pacific Journal of Mathematics*, **246**(2)(2010), 333-343.